

Mimicking the edge of 2d topological superconductors and insulators in 1d lattice models

Max Metlitski



Simons Collaboration on Ultra Quantum Matter: Kickoff Meeting.
Harvard, September 12, 2019

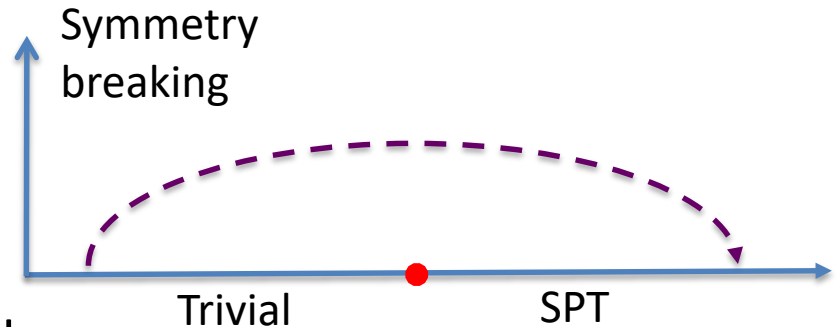


Robert Jones,
MIT

R. A. Jones and MM, arXiv:1902.05957
MM, arXiv:1908.08958

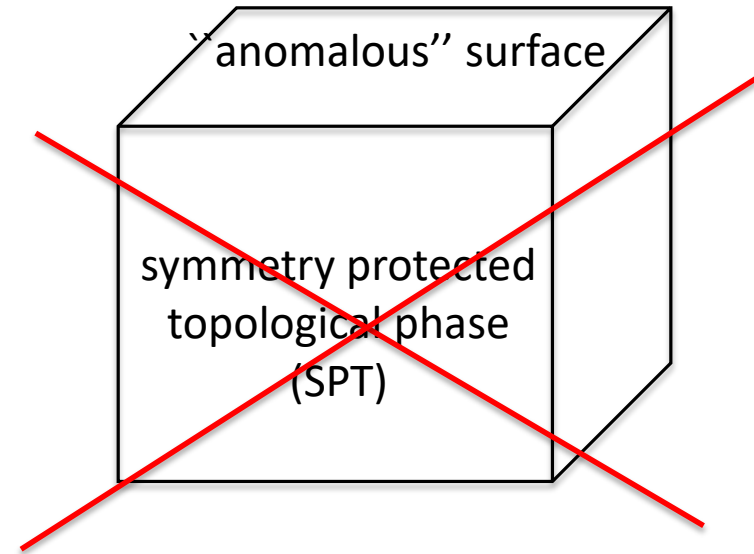
Symmetry protected topological phases (SPTs)

- Protected by symmetry G
- Gapped bulk
 - unique ground state on a closed spatial manifold
 - no anyon excitations



- Surface must be non-trivial
 - t'Hooft anomaly (can't couple to a background gauge-field).

Q: Can we realize/mimick the boundary in a lattice model without the bulk?



Motivation:

- i) Simulating edge of SPT without the bulk overhead
- ii) A more lattice based understanding of anomalies

Q: Can we realize/mimick the boundary in a lattice model without the bulk?

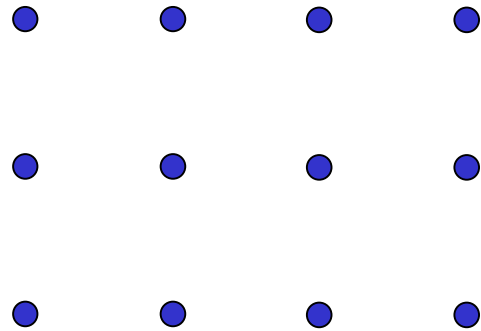
A: No, subject to

i) $V = \bigotimes_i V_i$

ii) $U(g) = \prod_i U_i(g)$

(On-site symmetry)

- Can always be gauged!



$$U_i(g)U_i(h) = U_i(gh)$$

Cohomology phases

i) $V = \otimes_i V_i$

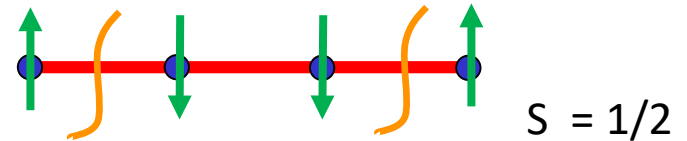
ii) ~~$U(g) = \prod_i U_i(g)$~~

- Class of boson SPT phases X. Chen, Z.-C. Gu, Z.-X. Liu, X.G.-Wen (2012)
- Classified by $H^{d+1}(G, U(1))$
(when G is gauged – Dijkgraaf-Witten theories)
- Edge can be mimicked if ii) is dropped (discrete G)
- 1+1D boundary: extract anomaly $H^3(G, U(1))$

Chen, Liu, Wen (2011); D. Else, C. Nayak (2014)

Example

- $G = \mathbb{Z}_2$
- $H^3(G, U(1)) = \mathbb{Z}_2$
- 1+1D boundary



$$U = (-1)^{N_{dw}/2} \prod_i \sigma_i^x$$

Cohomology phases

i) $V = \otimes_i V_i$

ii) ~~$U(g) = \prod_i U_i(g)$~~

- Extends to “super-cohomology” phases of fermions

(explicit construction for 1+1D boundary and discrete G in MM, arXiv:1908.08958)

This talk

- Beyond (super)cohomology phases:
 - 2+1D fermion topological superconductors

$$\cancel{V = \bigotimes_i V_i}$$

$$\cancel{U(g) = \prod_i U_i(g)}$$

This talk

- (Super)cohomology phases with a continuous symmetry group G
 - infinite dimensional site Hilbert space?
- 2+1D fermion topological insulators (quantum spin Hall)

$$V = \otimes_i V_i$$

~~$$U(g) = \prod_i U_i(g)$$~~

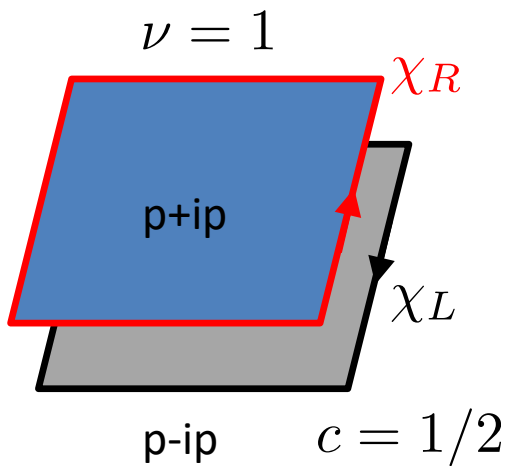
MM, arXiv:1908.08958

See also J. H. Son and J. Alicea, 1906.11846

Beyond (super)-cohomology

- Ex: 2+1D, fermions, $G = Z_2 \times Z_2^f$, $\nu \in Z_8$

Gu, Levin (14)



$$Z_2 : \chi_R \rightarrow -\chi_R, \quad \chi_L \rightarrow \chi_L$$

$$\delta L = im\chi_R\chi_L$$

$$m \rightarrow -m$$

Kramers-Wannier duality:

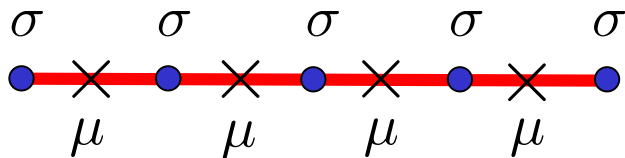
$$\text{High T} \longleftrightarrow \text{low T}$$

- Mimick in 1+1D?

Kramers-Wannier

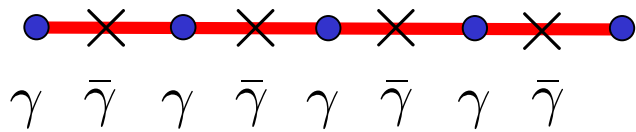
$$H_{TFIM} = -J \sum_i \sigma_i^x \sigma_{i+1}^x + h \sum_i \sigma_i^z$$

$$\mu_{i+1/2}^z = -\sigma_i^x \sigma_{i+1}^x, \quad \mu_{i-1/2}^x \mu_{i+1/2}^x = -\sigma_i^z \quad J \leftrightarrow h$$

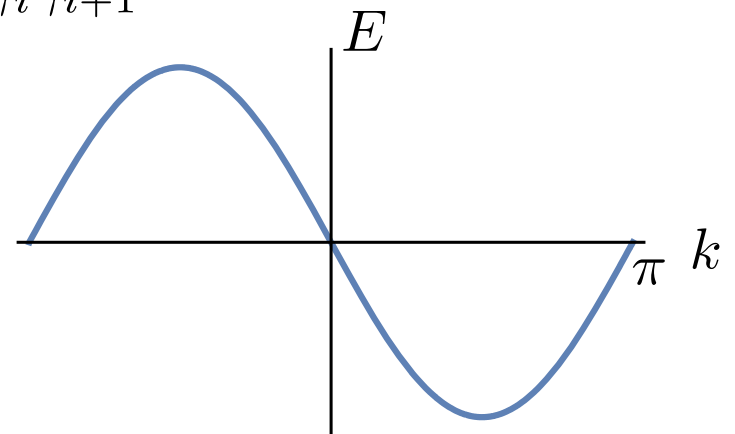


$$U_{KW}^2 = T_x$$

- Fermionize: $H = ih \sum_i \gamma_i \bar{\gamma}_i + iJ \sum_i \bar{\gamma}_i \gamma_{i+1}$



\rightarrow
 U_{KW}



$$\text{i) } V = \bigotimes_i V_i \qquad V_{phys} \subset V$$

$$\text{ii) } U(g) = \prod_i U_i(g)$$

- not enough to drop this if $U^2 = 1$

- Starting point:
Exactly solvable bulk model of Tarantino and Fidkowski (2016)

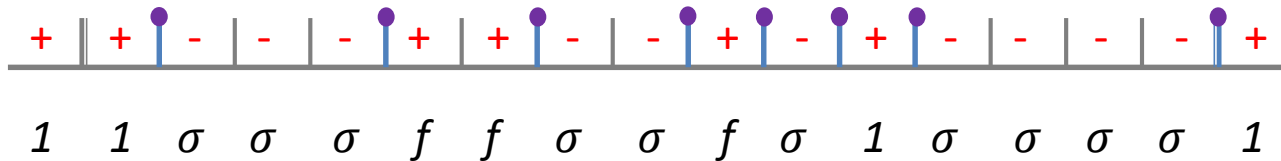
Intuition

$$\delta L = im\chi_R\chi_L$$



- Restore symmetry – Majorana liquid

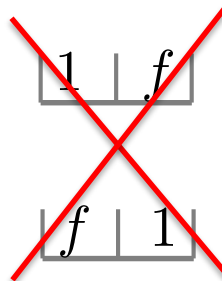
$$\dim V_{Major} = 2^{N_{dw}/2}$$



Link labels: $\{1, \sigma, f\}$,

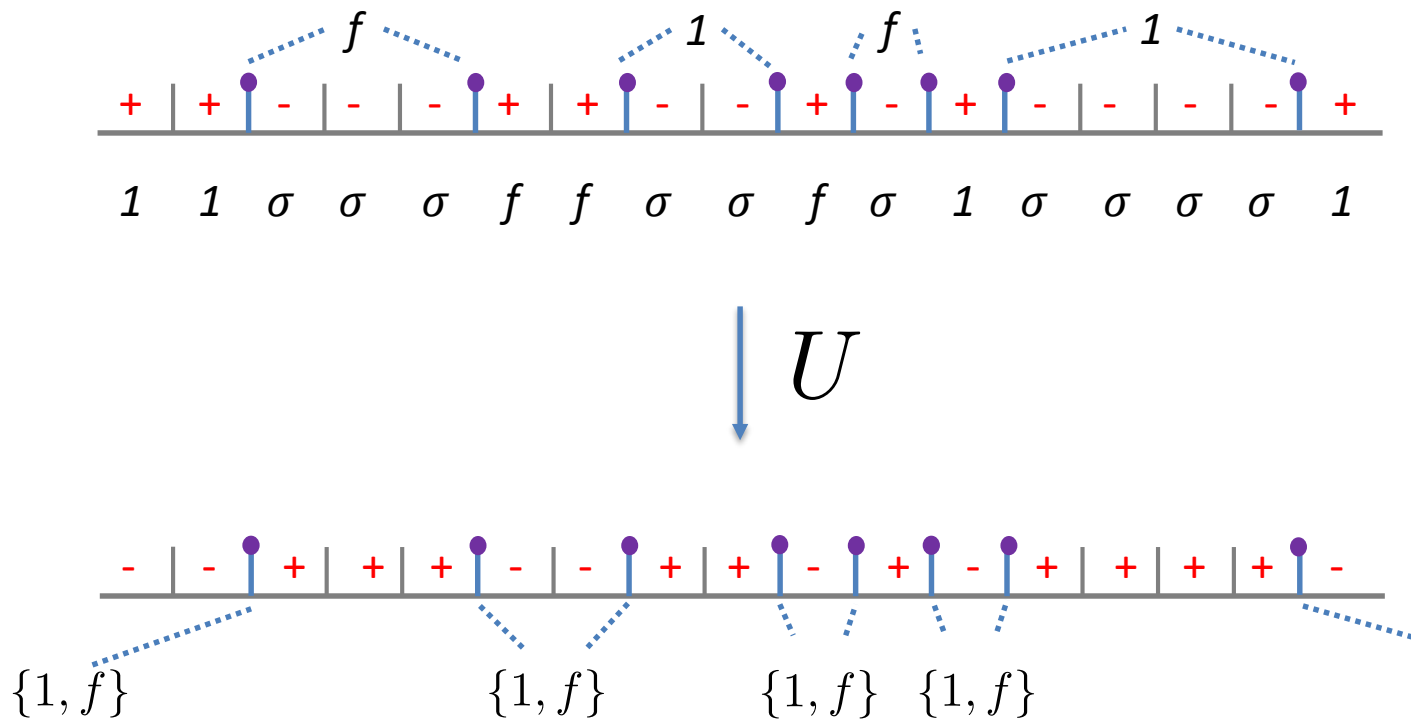
$$\sigma \times \sigma = 1 + f$$

- σ
- + 1 or f



$$\dim(V_{phys}) \sim (\sqrt{2} + 1)^N$$

Symmetry action



- Strings: $\{1, f\} \rightarrow \sigma$
 $\sigma \rightarrow \{1, f\}$ -superposition

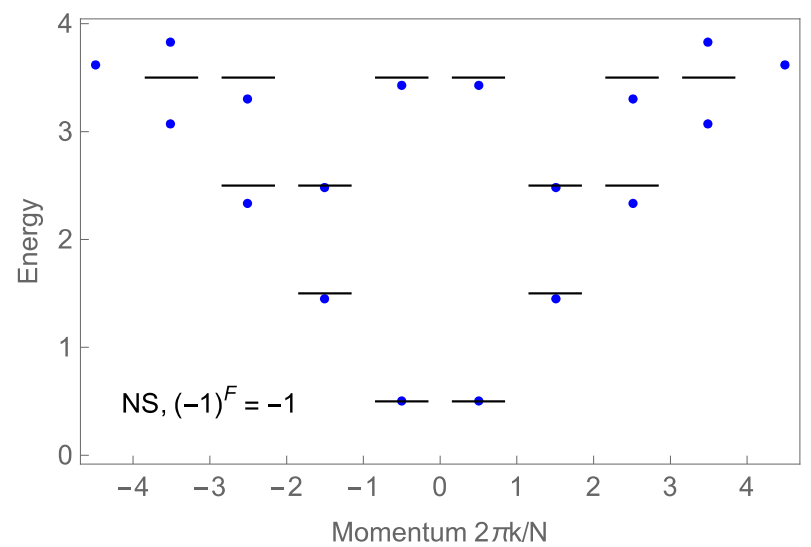
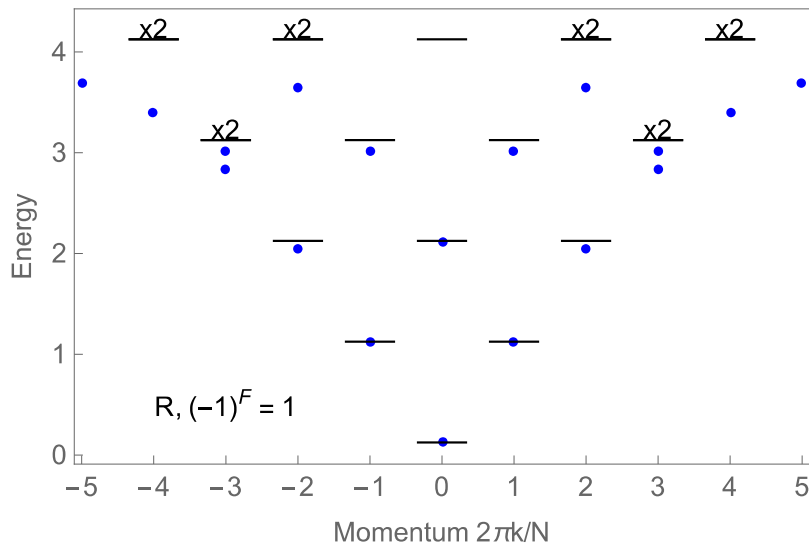
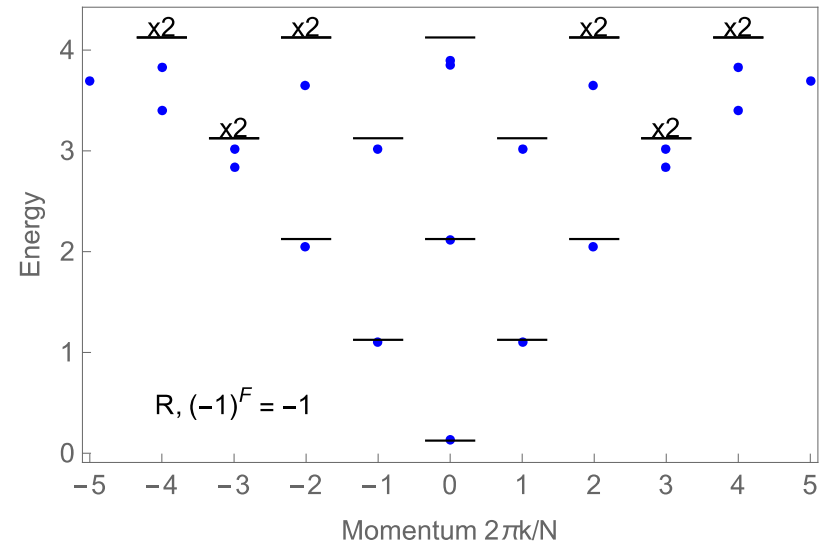
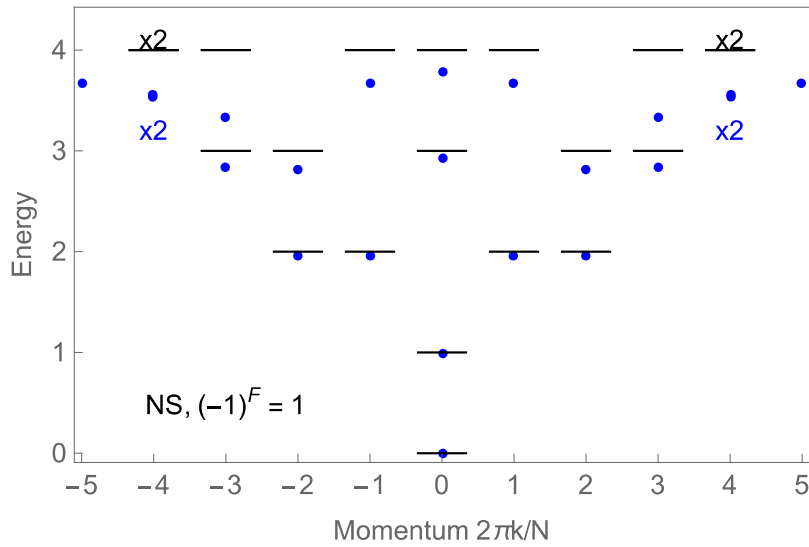
Hamiltonian



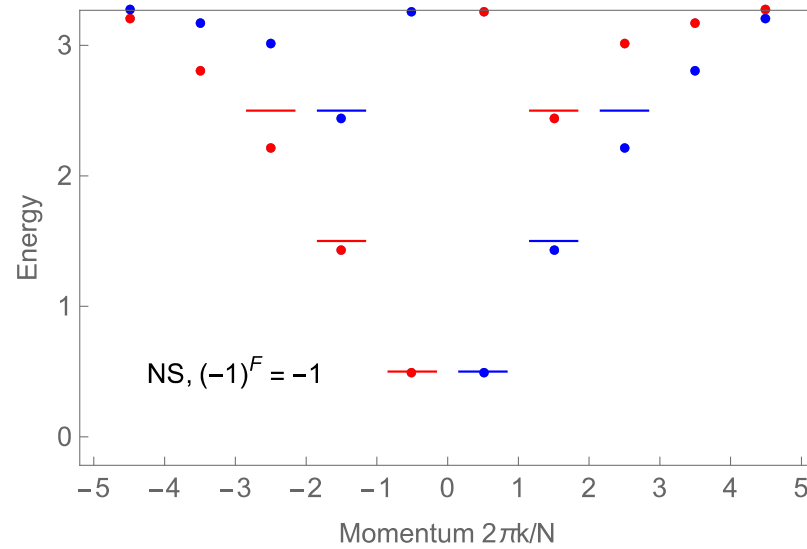
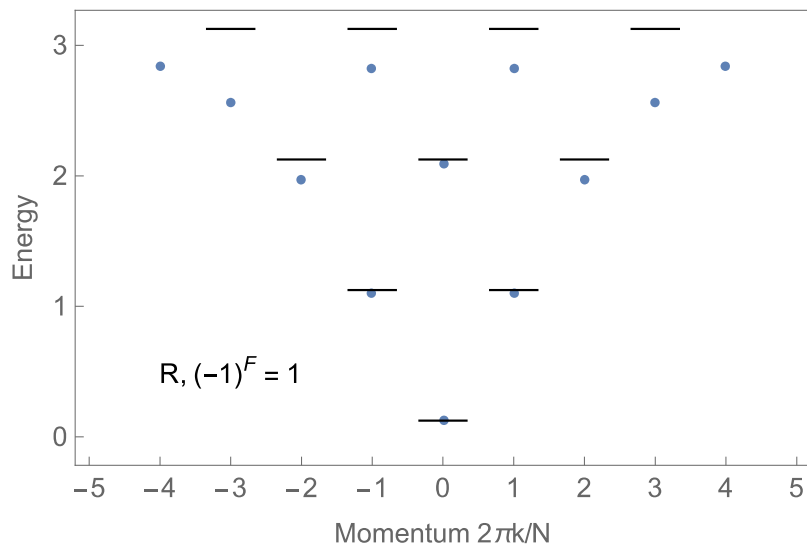
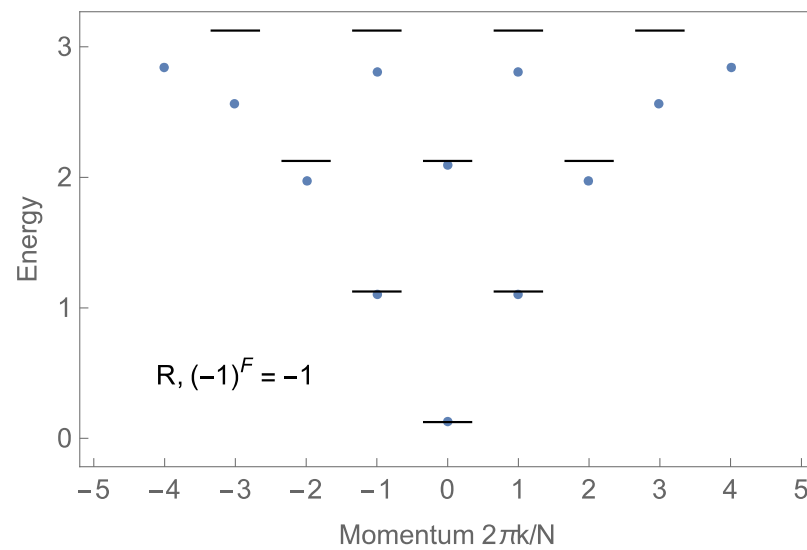
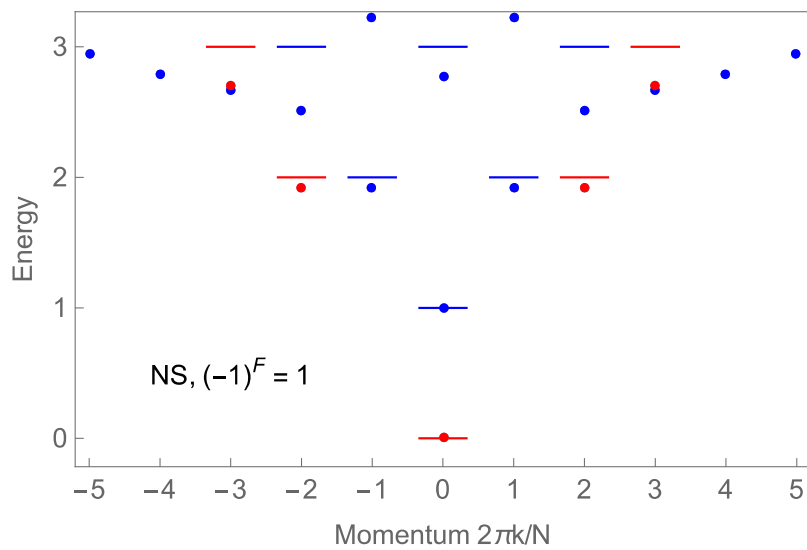
- $c_1 \approx c_2$ - solid numerical evidence (ED & DMRG) for Ising CFT

$$Z_2 : \chi_R \rightarrow -\chi_R, \quad \chi_L \rightarrow \chi_L$$

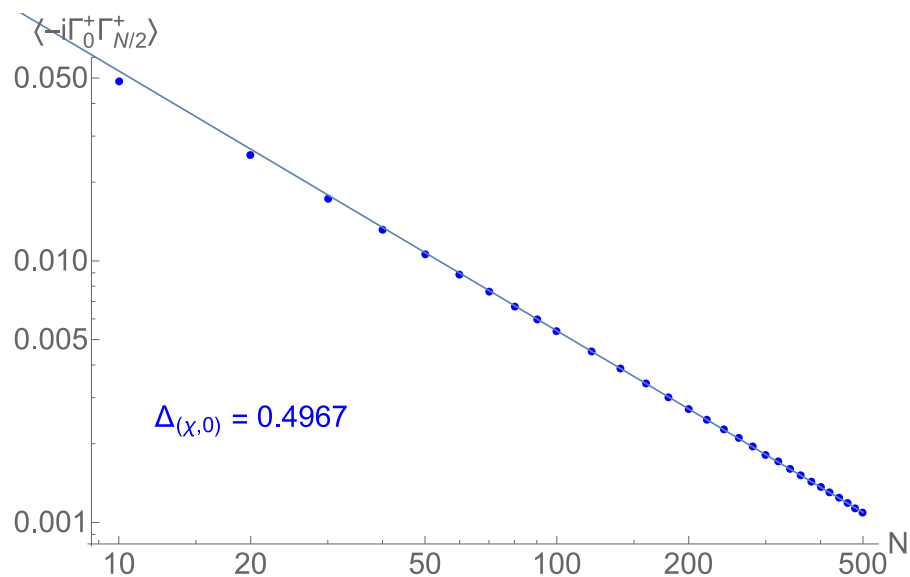
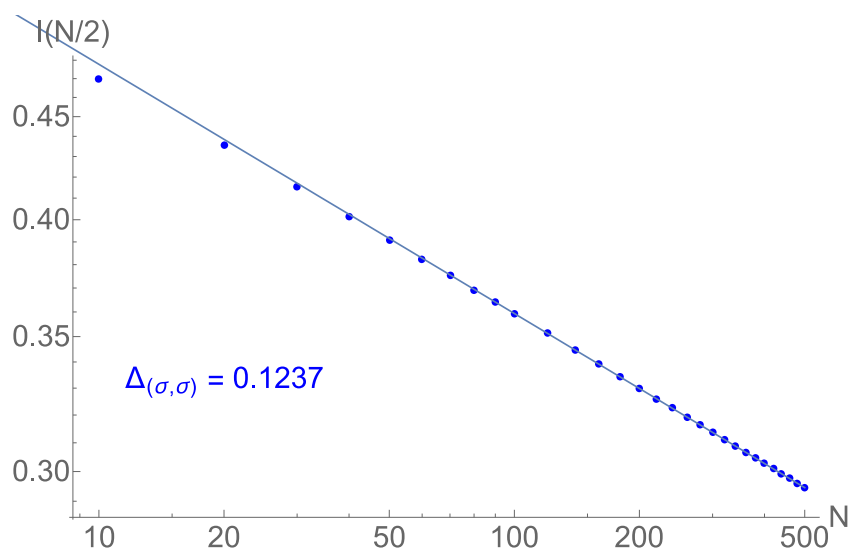
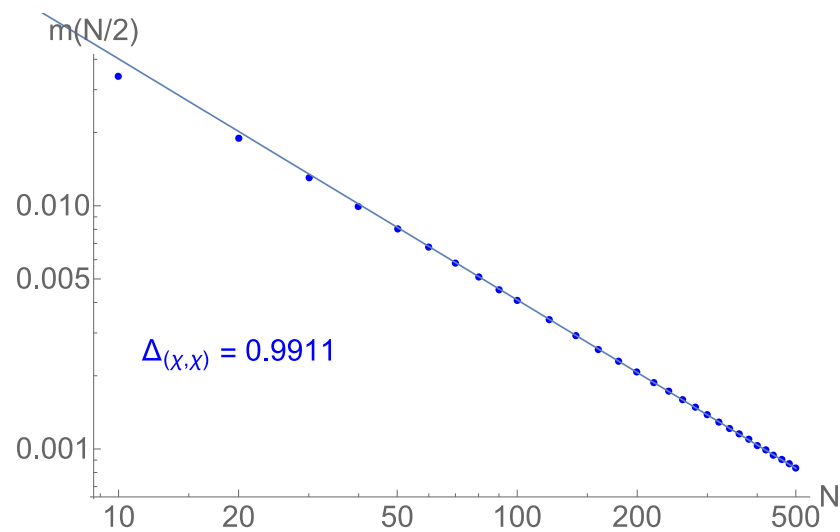
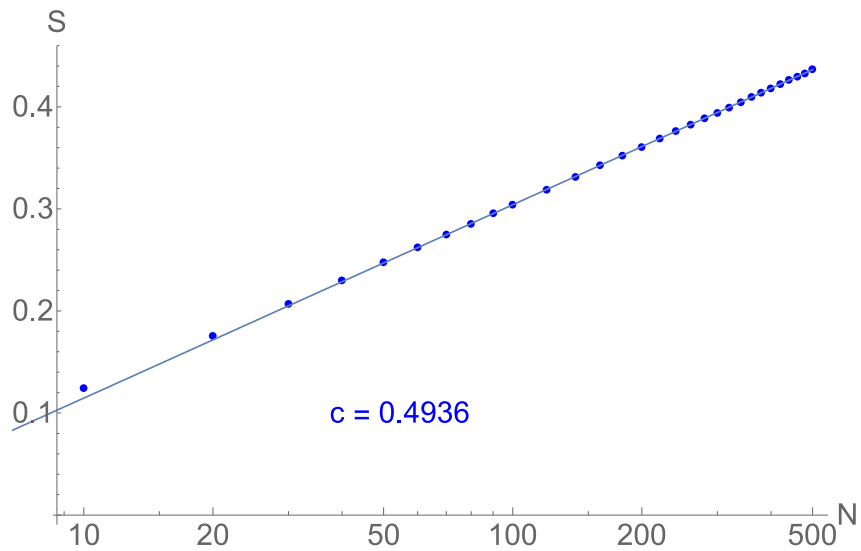
Exact diagonalization



Exact diagonalization

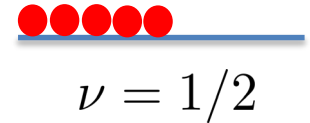


DMRG



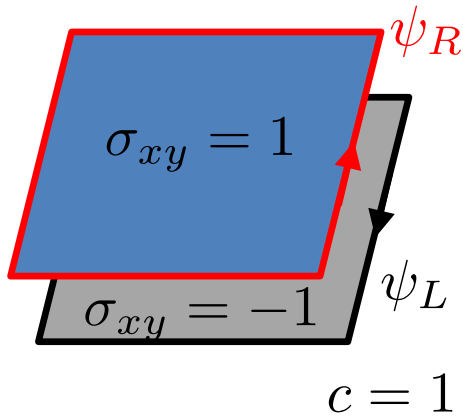
Lessons & open questions

- Mimick edge of 2+1D beyond super-cohomology fermion SPT
 - constrained Hilbert space
 - similar to 3+1D, U(1)xT fermion SPTs
 - generalizes to 2+1D fSPTs with $G \times Z_2^f$
 - obstruction?
 - other dimensions?



Quantum spin-Hall

- 2+1D, fermions




$$U : \psi_R \rightarrow -\psi_R, \quad \psi_L \rightarrow \psi_L$$

$$U(1) : N = N_R + N_L$$

$$\mathcal{T}_{NK} : \psi_R \leftrightarrow \psi_L, \quad i \rightarrow -i. \quad \mathcal{T}_{NK}^2 = 1$$

$$\mathcal{T} = U\mathcal{T}_{NK}, \quad \mathcal{T}^2 = (-1)^F$$

QSH edge

$$c_i = \frac{1}{2}(\gamma_i + i\bar{\gamma}_i)$$


$$g_{i,i+1} \in \{0, 1\}$$

$$\tau_i^z = (-1)^{g_i}$$

$$U = \left(\prod_{j=1}^L \tau_{j,j+1}^x \right) \left(\prod_{j=1}^L \gamma_j^{g_{j-1,j} + g_{j,j+1}} \right) (-i)^{N_{dw}/2}$$

- inspired by Tantisasadakarn and Vishwanath (18)

$$N = \sum_i n_i, \quad n_i = (-1)^{g_i(g_{i+1}+1)} c_i^\dagger c_i$$

$$\mathcal{T}_{NK} = \left(\prod_{j=1}^L (-1)^{g_{j,j+1}} c_j^\dagger c_j \right) \mathcal{T}_0$$

Properties

- \mathcal{T} domain wall:

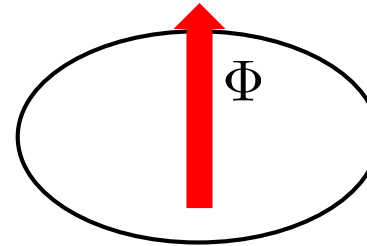


$$N \in 2\mathbb{Z} + 1$$

- Kramers parity switching

$$\Phi = 0, \quad \mathcal{T}^2 = (-1)^F$$

$$\Phi = \pi, \quad \mathcal{T}^2 = -(-1)^F$$



- Symmetric Hamiltonians
 - analyze by Jordan-Wigner
 - Luttinger liquid of QSH edge

- Can extract anomaly data: $G_b = G_f / Z_2^f$

$$\lambda \in H^2(G_b, Z_2), \quad \sigma \in H^2(G_b, Z_2), \quad w_3 \in C^3(G_b, Z_2)$$

$$d_T w_3 = (-1)^{\sigma \cup \sigma + \lambda \cup \sigma}$$

Open questions

- Mimick edge of 2+1D quantum spin-Hall
 - mimick other SPTs with continuous symmetry groups?
 - continuous vs discrete anomalies?

Thank you!